

Diracovo polje (jednacična)

$$\boxed{\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}}$$

Dirakove matrice $\{\gamma^\mu, \gamma^\nu\} = 2$

$$\gamma_0 = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix}_{4 \times 4}$$

$$\vec{\gamma} = \begin{pmatrix} 0 & \vec{\sigma} \\ -\vec{\sigma} & 0 \end{pmatrix}_{4 \times 4}$$

$$\vec{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$$

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\gamma^5 = i \gamma^0 \gamma^1 \gamma^2 \gamma^3$$

$$\gamma_5 = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

$$\gamma_5 = -i \gamma_0 \gamma_1 \gamma_2 \gamma_3$$

$$\gamma_{\mu\nu} = \frac{i}{2} [\gamma_\mu, \gamma_\nu] \Rightarrow \gamma^{\mu\nu} = \frac{i}{2} [\gamma^\mu, \gamma^\nu]$$

$$d = a^\mu \gamma_\mu = \cancel{a^\mu \gamma_\mu} a_\mu \gamma^\mu$$

$$\underline{(\gamma^0)^2 = 1} \quad \underline{(\gamma^i)^2 = -1}$$

Representacija je matrica i je γ^5 matrica

$$\gamma^5 = \begin{pmatrix} -\mathbb{I} & 0 \\ 0 & \mathbb{I} \end{pmatrix}_* , \quad \gamma^5 = \begin{pmatrix} 0 & \mathbb{I} \\ \mathbb{I} & 0 \end{pmatrix}_D \quad (\text{tr } \gamma^5 = 0)$$

$$\gamma^5 = -\frac{i}{4!} \epsilon^{\mu\nu\rho\sigma} \gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma$$

$$\gamma^0 = \begin{pmatrix} 0 & \mathbb{I} \\ \mathbb{I} & 0 \end{pmatrix}_* \quad \gamma^i = \begin{pmatrix} 0 & \gamma^i \\ -\gamma^i & 0 \end{pmatrix}_* \quad \left. \begin{array}{l} * - \text{hiral} \\ \text{na represen} \\ \text{cija} \end{array} \right\}$$

$$\gamma^0 = \begin{pmatrix} \mathbb{I} & 0 \\ 0 & -\mathbb{I} \end{pmatrix}_D \quad \gamma^k = \begin{pmatrix} 0 & \gamma^k \\ -\gamma^k & 0 \end{pmatrix}_D \quad \left. \begin{array}{l} D - \text{Dirac} \\ \text{representa} \end{array} \right\}$$

$$\begin{aligned} \gamma^5 &= i \gamma^0 \gamma^1 \gamma^2 \gamma^3 \\ &= i g^{00} \gamma_0 \gamma^1 \gamma_2 \gamma^2 \gamma_3 \gamma^3 \\ &= i \underbrace{g^{00} g^{11} g^{22} g^{33}}_{(-1)} \gamma_0 \gamma_1 \gamma_2 \gamma_3 \\ &= -i \gamma_0 \gamma_1 \gamma_2 \gamma_3 = \gamma_5 \end{aligned}$$

Hiralna reprezentacija ili Weyl-ova reprezentacija

$$(\gamma^5 = \gamma_5)$$

$$\gamma^5 = -\frac{i}{4!} \epsilon^{\mu\nu\rho\sigma} \gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma =$$

$$= -\frac{i}{4!} \epsilon^{\mu\nu\rho\sigma} g_{\mu\mu'} \gamma^{\mu'} g_{\nu\nu'} \gamma^{\nu'} g_{\rho\rho'} \gamma^{\rho'} g_{\sigma\sigma'} \gamma^{\sigma'}$$

$$= -\frac{i}{4!} \epsilon_{\mu\nu\rho\sigma} \gamma^{\mu'} \gamma^{\nu'} \gamma^{\rho'} \gamma^{\sigma'}$$

$$L_i = \begin{pmatrix} 0 & \gamma^i \\ \gamma^i & 0 \end{pmatrix}_D \rightarrow \text{Diracova reprezentacija}$$

$$\not{P} = i \not{P}$$

$$\not{P} = \gamma^\mu P_\mu = g_{\mu\nu} \gamma^\mu P^\nu = \gamma^0 P^0 - \vec{\gamma} \cdot \vec{P}$$

$$\not{P} = \gamma^0 P_0 - \gamma^i P_i$$

$$P_0 \equiv \frac{E}{c} = i \frac{\partial}{\partial x^0}, \quad P_i = -i \frac{\partial}{\partial x^i} = -i \partial_i$$

$$\begin{aligned} \not{P} &= i \gamma^0 \frac{\partial}{\partial x^0} + i \gamma^i \partial_i \equiv i \gamma^0 \partial_0 + i \gamma^i \partial_i \\ &= i \gamma^\mu \partial_\mu = i \not{P} \end{aligned}$$

$$\not{P} = i \not{P}$$

$$\not{P} = \gamma^0 \partial_0 + \gamma^i \partial_i$$

1. Nova je data Dirakova $\bar{\psi}$ -na

$$i\hbar \frac{\partial \psi}{\partial t} = [-i\hbar c \vec{\alpha} \cdot \vec{\nabla} + \beta m c^2] \psi$$

gde su $\vec{\alpha}$ i β nepoznate. Da bi ova $\bar{\psi}$ -na opisivala slobodnu relativističnu česticu, potrebno je da zadovoljava KG $\bar{\psi}$ -na. Na osnovi ovog ograničenja, nađi $\vec{\alpha}$ i β .

Primeniti operator

$$i \partial_t = (-i \vec{\alpha} \cdot \vec{\nabla} + \beta m) \text{ na levu i desnu stranu Dirakove } \bar{\psi}\text{-ne. } (c = \hbar = 1)$$

Sledi

$$-\frac{\partial^2 \psi}{\partial t^2} = - \left(\alpha_i \frac{\partial}{\partial x_i} \right) \left(\alpha_j \frac{\partial}{\partial x_j} \right) \psi - i m \{ \alpha_i, \beta \} \frac{\partial \psi}{\partial x_i} + m^2 \beta^2 \psi \quad (1)$$

KG $\bar{\psi}$ -na

$$(\partial^\mu \partial_\mu + m^2) \psi = 0$$

$$\frac{\partial^2 \psi}{\partial t^2} - \Delta \psi + m^2 \psi = 0$$

$$\frac{\partial^2 \psi}{\partial t^2} = + \Delta \psi + m^2 \psi$$

↓

$$-\frac{\partial^2 \psi}{\partial t^2} = -\Delta \psi + m^2 \psi$$

$$-\alpha_i^2 \frac{\partial^2 \psi}{\partial x_i^2} - (\alpha_i \alpha_j + \alpha_j \alpha_i) \frac{\partial^2 \psi}{\partial x_i \partial x_j}$$

(2)

Poređenjem (1) i (2)

$$d_i d_j + d_j d_i = 2\delta_{ij}$$

$$\{d_i, \beta\} = 0$$

(3)

$$\beta^2 = 1$$

Ako su d_i i β brojevi, onda će prve dve relacije u (3) biti zadovoljene samo ako su bar 3 od 4 broja nula. A to se ne slaže sa $d_i^2 = \beta^2 = 1$. Dakle, u pitanju su ^{matrice!}

Ako se uvede oznaka $1_0 = \beta$, onda se

(3) može zapisati kao

$$\{d_\mu, d_\nu\} = 2\delta_{\mu\nu}$$

Ove matrice moraju biti Hermitske.

Takođe, one su tražena nula

$$d_\mu d_\nu = -d_\nu d_\mu$$

$$d_\mu^2 d_\nu = -d_\nu d_\mu d_\nu$$

$$\text{Tr}(d_\nu) = -\text{Tr}(d_\mu d_\nu d_\mu) = -\text{Tr}(d_\mu^2 d_\nu) = -\text{Tr}(d_\nu)$$

Zato što je $d_\mu^2 = 1$, svojstvene vrednosti su ± 1 . Dakle d_μ su parne dimenzije ($2 \times 2, 4 \times 4$) npr.

U slučaju 2x2 matrice, imamo Paulyjeve matrice koje zadovoljavaju $\{\beta_i, \beta_j\} = 2\delta_{ij}$.
 Sa jediničnom matricom čine bazu u \mathbb{C}^{22} .
 Ali, $\{\beta_i, I\} = 2\beta_i$ umesto nulte Dake, mora se ići na 4x4 matrice. Sledeći izbor se pokazuje kao takav:

$$d_0 = \begin{pmatrix} \mathbf{I} & 0 \\ 0 & -\mathbf{I} \end{pmatrix}_{4 \times 4}, \quad d_i = \begin{pmatrix} 0 & \beta_i \\ \beta_i & 0 \end{pmatrix}_{4 \times 4}$$

↑
Pali znak? (H)

Drugi zapis

$$j^0 = \beta \quad j^i = \beta d_i$$

$$i\hbar j^0 \frac{\partial \psi}{\partial t} = (-i\hbar c j^i \nabla_i + mc^2) \psi / c$$

$$i j^0 \frac{\partial \psi}{\partial x^0} = (-i j^i \frac{\partial}{\partial x^i} + m) \psi \quad ; t=1, c=1$$

$$i j^m \frac{\partial \psi}{\partial x^m} - m\psi = 0$$

→
v. poradu

$$(i j^m \partial_m - m) \psi$$

$$(i \not{\partial} - m) \psi = 0$$

$$\begin{aligned} 0 & \rightarrow j^m A_m = \cancel{A} \\ 1 & \rightarrow j^0 A^0 - \vec{j} \cdot \vec{A} = \cancel{A} \end{aligned}$$

$$i\hbar \frac{\partial \psi}{\partial t} = H_D \psi$$

$$H_D = c \vec{\alpha} \cdot \vec{p} + \beta E_0$$

→
B. poradu

2. Pokazati da iz $\{k_\mu, k_\nu\} = 2g_{\mu\nu}$ sledi

$\{b_i, b_j\} = 2\delta_{ij}$. Koristiti det:

$$k_0 = B \quad i \quad k_i = B d_i, \quad d_i = \begin{pmatrix} 0 & b_i \\ b_i & 0 \end{pmatrix}$$

Bitno za nerelativistički limit.

$$\{k_i, k_j\} = 2g_{ij} I = -2\delta_{ij} I$$

$$\{B d_i, B d_j\} = -2\delta_{ij} I$$

$$B d_i B d_j + B d_j B d_i = -2\delta_{ij} I$$

(pa se koristi
anti-komutator
 $\{B, d_i\} = 0$)

$$-d_i d_j - d_j d_i = -2\delta_{ij} I$$

$$\boxed{\{d_i, d_j\} = 2\delta_{ij} I}$$

$$\begin{pmatrix} 0 & b_i \\ b_i & 0 \end{pmatrix} \begin{pmatrix} 0 & b_j \\ b_j & 0 \end{pmatrix} + \begin{pmatrix} 0 & b_j \\ b_j & 0 \end{pmatrix} \begin{pmatrix} 0 & b_i \\ b_i & 0 \end{pmatrix} = 2\delta_{ij} I$$

$$\begin{pmatrix} b_i b_j & 0 \\ 0 & b_i b_j \end{pmatrix} + \begin{pmatrix} b_j b_i & 0 \\ 0 & b_j b_i \end{pmatrix} = 2\delta_{ij} I$$

$$\begin{pmatrix} \{b_i, b_j\} & 0 \\ 0 & \{b_i, b_j\} \end{pmatrix} = \begin{pmatrix} 2\delta_{ij} & 0 \\ 0 & 2\delta_{ij} \end{pmatrix}$$

$$\Downarrow$$
$$\boxed{\{b_i, b_j\} = 2\delta_{ij}}$$

3. Nena je $U = e^{\beta \vec{L} \cdot \vec{n}}$, gde su β i \vec{L} Diranove matrice, \vec{n} je jedinični vektor. Pokazati da važi:

$$\vec{L}' = U \vec{L} U^\dagger = \vec{L} - (1 - U^2) (\vec{L} \cdot \vec{n}) \vec{n}$$

BHC f-la

$$e^B e^A = e^{A + [B, A] + \frac{1}{2!} [B, [B, A]] + \dots}$$

$$\left[\beta L_i = -L_i \beta \quad \beta^\dagger = \beta, \quad L_i^\dagger = L_i \right]$$

$$U = e^{\beta \vec{L} \cdot \vec{n}}, \quad U^\dagger = ?$$

$$\begin{aligned} U^\dagger &= \left(e^{\beta \vec{L} \cdot \vec{n}} \right)^\dagger = \left(e^{\beta L_i n_i} \right)^\dagger = e^{(\beta L_i)^\dagger n_i} \\ &= e^{L_i \beta^\dagger n_i} = e^{L_i \beta n_i} = e^{-\beta L_i n_i} = e^{-\beta \vec{L} \cdot \vec{n}} \end{aligned}$$

$$U^\dagger = e^{-\beta \vec{L} \cdot \vec{n}}$$

$$\vec{L}' = e^{\beta \vec{L} \cdot \vec{n}} \vec{L} e^{-\beta \vec{L} \cdot \vec{n}}$$

$$= \vec{L} + [\beta \vec{L} \cdot \vec{n}, \vec{L}] + \frac{1}{2!} [\beta \vec{L} \cdot \vec{n}, [\beta \vec{L} \cdot \vec{n}, \vec{L}]] + \dots$$

$$\begin{aligned}
 [\beta \vec{l} \cdot \vec{n}, \vec{L}] &= [\beta l_j n_j, L_i] = n_j [\beta l_j, L_i] \\
 &= n_j (\beta l_j L_i - L_i \beta l_j) = n_j (\beta l_j L_i + \beta L_i l_j) \\
 &= n_j \beta \{l_j, L_i\} = 2\delta_{ij} n_j \beta = 2\beta n_i = 2\beta \vec{n}
 \end{aligned}$$

$$\begin{aligned}
 [\beta l_i n_i, 2\beta n_j] &= 2n_i n_j [\beta l_i, \beta] \\
 &= 2n_i n_j (\beta l_i \beta - \beta^2 l_i) = 2n_i n_j (-l_i - l_i) \\
 &= -4n_i n_j l_i = -4\vec{n} (\vec{L} \cdot \vec{n})
 \end{aligned}$$

$$\begin{aligned}
 [\beta l_i n_i, \underbrace{-4n_i n_j}_{\delta_{ij}} l_i] &= -4n_i \delta_{ij} (n_j l_i l_i - l_i \beta l_i) \\
 &= -4n_i \delta_{ij} (\beta l_i^2 + \beta l_i^2) = -8n_i \delta_{ij} \beta = -8\vec{n} \beta
 \end{aligned}$$

$$\begin{aligned}
 [\beta l_i n_i, -8n_j \beta] &= -8n_i n_j (\beta l_i \beta - \beta^2 l_i) \\
 &= -8n_i n_j (-l_i - l_i) \\
 &= 16n_i n_j l_i = 16\vec{n} (\vec{L} \cdot \vec{n})
 \end{aligned}$$

$$\vec{L}' = \vec{L} + 2\beta \vec{n} - 2\vec{n} (\vec{L} \cdot \vec{n}) - \frac{8}{3!} \vec{n} \beta + \frac{16}{4!} \vec{n} (\vec{L} \cdot \vec{n})$$

Desna strana izrazu

$$\vec{L}' = \vec{L} - (1 - v^2) (\vec{L} \cdot \vec{u}) \vec{u} \quad ; \quad v = e^{\beta \vec{L} \cdot \vec{u}}$$

$$v^2 = e^{2\beta \vec{L} \cdot \vec{u}} = 1 + 2\beta \vec{L} \cdot \vec{u} + \frac{(2\beta \vec{L} \cdot \vec{u})^2}{2!} +$$

$$+ \frac{(2\beta \vec{L} \cdot \vec{u})^3}{3!} + \frac{(2\beta \vec{L} \cdot \vec{u})^4}{4!} + \dots$$

$$\begin{aligned} (\beta \vec{L} \cdot \vec{u})^2 &= \beta^2 \delta_{ik} \delta_{jl} u_i u_j = \delta_{ik} \delta_{jl} \beta^2 u_i u_j \\ &= -\delta_{ik} \delta_{jl} u_i u_j = -\delta_{ij} u_i u_j = -u_i^2 = -1 \end{aligned}$$

$$\begin{aligned} (\beta \vec{L} \cdot \vec{u})^3 &= (\beta \vec{L} \cdot \vec{u}) (\beta \vec{L} \cdot \vec{u})^2 = -(\beta \vec{L} \cdot \vec{u}) \\ &= -\beta \vec{L} \cdot \vec{u} \end{aligned}$$

$$(\beta \vec{L} \cdot \vec{u})^4 = 1$$

$$v^2 = 1 + 2\beta \vec{L} \cdot \vec{u} + \frac{4}{2} (-1) - \frac{8}{3!} \beta \vec{L} \cdot \vec{u} + \frac{16}{4!} + \dots$$

$$v^2 = 1 + 2\beta \vec{L} \cdot \vec{u} - 2 - \frac{8}{3!} \beta \vec{L} \cdot \vec{u} + \frac{16}{4!} + \dots$$

$$1 - v^2 = 4 - 4 - 2\beta \vec{L} \cdot \vec{u} + 2 + \frac{8}{3!} \beta \vec{L} \cdot \vec{u} - \frac{16}{4!} + \dots$$

$$-(1 - v^2) = 2\beta \vec{L} \cdot \vec{u} - 2 - \frac{8}{3!} \beta \vec{L} \cdot \vec{u} + \frac{16}{4!}$$

$$\vec{L}' = \vec{L} + \left(2\beta \vec{L} \cdot \vec{n} - 2 \frac{\beta}{3!} \beta \vec{L} \cdot \vec{n} + \frac{16}{4!} \dots \right) (\vec{n} \cdot \vec{L}) \vec{n}$$

$$\vec{L}' = \vec{L} + 2\beta (\vec{L} \cdot \vec{n}) \vec{n} - 2(\vec{n} \cdot \vec{L}) \vec{n} - \frac{8}{3} \beta (\vec{n} \cdot \vec{L})^2 \vec{n} + \frac{16}{4!} (\vec{n} \cdot \vec{L}) \vec{n} + \dots$$

$$\vec{L}' = \vec{L} + 2\beta \vec{n} - 2\vec{n} (\vec{L} \cdot \vec{n}) - \frac{8}{3} \vec{n} \beta + \frac{16}{4!} (\vec{n} \cdot \vec{L}) \vec{n} + \dots$$

što je jednako levoj strani izrazu iz postavke.

$$|\vec{n} \cdot \vec{L}'|^2 = \sum_i n_i'^2 L_i'^2 = \sum_i \delta_{ij} L_i'^2 L_j'^2 = L_i'^2 = 1 \quad \perp$$

↓
Važno!

4. Pokazati da je

$$1. \mu^{\mu+} = \mu^0 \mu^{\mu} \mu^0.$$

$$2. \delta^{\mu\nu+} = \mu^0 \delta^{\mu\nu} \mu^0$$

$$1. (\mu^0)^+ = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix}^+ = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix} = \mu^0 \mu^0 \mu^0 = \mu^0$$

$$(\mu^i)^+ = \begin{pmatrix} 0 & \delta_{ij} \\ -\delta_{ij} & 0 \end{pmatrix}^+ = - \begin{pmatrix} 0 & \delta_{ij} \\ -\delta_{ij} & 0 \end{pmatrix} =$$

$$= -(\mu^0)^2 \mu^i = -\mu^0 \mu^0 \mu^i = \mu^0 \mu^i \mu^0$$

$$(\mu^0)^2 = 1, \quad \{\mu^0, \mu^i\} = 0$$

$$2. \delta^{\mu\nu+} = -\frac{i}{2} (\mu^{\mu} \mu^{\nu} - \mu^{\nu} \mu^{\mu})^+$$

$$= -\frac{i}{2} (\mu^{\nu+} \mu^{\mu+} - \mu^{\mu+} \mu^{\nu+})$$

$$= -\frac{i}{2} (\mu^0 \mu^{\nu} \mu^0 \mu^0 \mu^{\mu} \mu^0 - \mu^0 \mu^{\mu} \mu^0 \mu^0 \mu^{\nu} \mu^0)$$

$$= -\frac{i}{2} (\underbrace{\mu^0 \mu^{\nu} \mu^{\mu} \mu^0}_{\text{blue}} - \underbrace{\mu^0 \mu^{\mu} \mu^{\nu} \mu^0}_{\text{blue}})$$

$$= -\frac{i}{2} \mu^0 [\mu^{\nu}, \mu^{\mu}] \mu^0$$

$$= \delta^{\mu\nu} \mu^0$$

5. Pokazati da je

$$a) \quad \mu_5^\dagger = \mu_5, \quad \mu_5^5 = \mu_5^{-1}$$

$$b) \quad (\mu_5)^2 = 1$$

$$c) \quad \{\mu_5, \mu_n\} = 0$$

$$d) \quad (\mu_5 \mu_n)^\dagger = \mu_0 \mu_5 \mu_n \mu_0$$

a) P. definiciji

$$\mu_5 = -i \mu_0 \mu_1 \mu_2 \mu_3$$

$$\mu_5^\dagger = i \mu_3^\dagger \mu_2^\dagger \mu_1^\dagger \mu_0^\dagger$$

$$= i \underbrace{\mu_0 \mu_3 \mu_0}_{1} \underbrace{\mu_0 \mu_2 \mu_0}_{1} \underbrace{\mu_0 \mu_1 \mu_0}_{1} \mu_0$$

$$= i \mu_0 \mu_3 \mu_2 \mu_1$$

$$= -i \mu_0 \mu_3 \mu_1 \mu_2 = i \mu_0 \mu_1 \mu_3 \mu_2$$

$$= -i \mu_0 \mu_1 \mu_2 \mu_3 = \mu_5$$

$$\mu_5^5 = \mu_5^{-1}$$

$$\mu_5^{-1} = (-i \mu_0 \mu_1 \mu_2 \mu_3)^{-1} = i \mu_3^{-1} \mu_2^{-1} \mu_1^{-1} \mu_0^{-1}$$

Poznato da je

$$\{ \mu_n, \mu^m \} = 2\delta_{nm}$$

$$\mu_0 \mu^0 = 1 \Rightarrow \mu_0^{-1} = \mu^0$$

$$\mu_i \mu^i = 1 \Rightarrow \mu_i^{-1} = \mu^i$$

$$\begin{aligned} \mu_5^{-1} &= i \mu^3 \mu^2 \mu^1 \mu^0 \\ &= -i \mu^3 \mu^2 \mu^0 \mu^1 = i \mu^3 \mu^0 \mu^2 \mu^1 = \\ &= -i \mu^0 \mu^3 \mu^2 \mu^1 = i \mu^0 \mu^2 \mu^3 \mu^1 = \\ &= -i \mu^0 \mu^2 \mu^1 \mu^3 = i \mu^0 \mu^1 \mu^2 \mu^3 = \mu_5 \end{aligned}$$

Sa druge strane:

$$\begin{aligned} \mu_5 &= i \mu^0 \mu^1 \mu^2 \mu^3 = i \mu_0 g^{00} \mu_1 g^{11} \mu_2 g^{22} \mu_3 g^{33} \\ &= i g^{00} g^{11} g^{22} g^{33} \mu_0 \mu_1 \mu_2 \mu_3 \\ &= -i \mu_0 \mu_1 \mu_2 \mu_3 = \mu_5 \end{aligned}$$

Dakle

$$\mu_5^{-1} = \mu_5 \Rightarrow (\mu_5)^2 = 1 \quad (b)$$

$$c) \{ \mu_5, \mu_6 \} = 0$$

↓
antikomutator

$$\begin{aligned}
\gamma_0 \gamma_5 &= -i \gamma_0 \gamma_1 \gamma_2 \gamma_3 \\
&= i \gamma_0 \gamma_4 \gamma_0 \gamma_2 \gamma_3 \\
&= -i \gamma_0 \gamma_2 \gamma_2 \gamma_0 \gamma_3 \\
&= i \gamma_0 \gamma_1 \gamma_2 \gamma_3 \gamma_0 \\
&= -(-i \gamma_0 \gamma_1 \gamma_2 \gamma_3) \gamma_0 \\
&= -\gamma_5 \gamma_0
\end{aligned}$$

↓

$$\{\gamma_0, \gamma_5\} = 0$$

$$\begin{aligned}
\gamma_1 \gamma_5 &= -i \gamma_1 \gamma_0 \gamma_1 \gamma_2 \gamma_3 \\
&= i \gamma_0 \gamma_1 \gamma_1 \gamma_2 \gamma_3 \\
&= -i \gamma_0 \gamma_1 \gamma_2 \gamma_1 \gamma_3 \\
&= i \gamma_0 \gamma_1 \gamma_2 \gamma_3 \gamma_1 \\
&= -(-i \gamma_0 \gamma_1 \gamma_2 \gamma_3) \gamma_1 \\
&= -\gamma_5 \gamma_1
\end{aligned}$$

$$\{\gamma_1, \gamma_5\} = 0$$

Što čemo ovomiz $\{\gamma_2, \gamma_5\} = 0 = \{\gamma_3, \gamma_5\} \Rightarrow$

$$\Rightarrow \{\gamma_\mu, \gamma_5\} = 0$$

$$d) (\gamma_5 \gamma_\mu)^+ = \gamma_\mu^+ \gamma_5^+ = \gamma_\mu^+ \gamma_5$$

$$= \gamma_0 \gamma_\mu \gamma_0 \gamma_5 = -\gamma_0 \gamma_\mu \gamma_5 \gamma_0 = \gamma_0 \gamma_5 \gamma_\mu \gamma_0$$

$$\boxed{(\gamma_5 \gamma_\mu)^+ = \gamma_0 \gamma_5 \gamma_\mu \gamma_0}$$

6. Pomažati da je

$$\mu_5 = -\frac{i}{4!} \epsilon_{\mu_1 \mu_2 \mu_3 \mu_4} \mu^{\mu_1} \mu^{\mu_2} \mu^{\mu_3} \mu^{\mu_4}$$

$$\epsilon_{\mu_1 \mu_2 \mu_3 \mu_4} = \begin{cases} +1 & \text{Za parne permutacije} \\ & \text{od } 0123 \\ -1 & \text{Za neparne per. od } 0123 \\ 0 & \text{ako su dva indeksa ista} \end{cases}$$

Permutacije od 0123

- 0123	- 1023	- 2013	- 3012
- 0132	+ 1032	+ 2031	+ 3021
- 0312	- 1302	- 2301	- 3201
- 0321	+ 1320	+ 2310	+ 3210
- 0213	- 1230	- 2130	- 3120
- 0231	- 1203	+ 2103	+ 3102

Ukupno 24 permutacija

Dakle na osnovi gorejih permutacija
biće na primer

$$\mu_5 = -\frac{i}{4!} \left(\mu^0 \mu^1 \mu^2 \mu^3 - \mu^0 \mu^1 \mu^3 \mu^2 + \mu^0 \mu^3 \mu^1 \mu^2 - \mu^0 \mu^3 \mu^2 \mu^1 + \mu^0 \mu^2 \mu^1 \mu^3 - \mu^0 \mu^2 \mu^3 \mu^1 + \dots \text{itd.} \right)$$

Konšidenjam pelacov

$$\{ \mu^4, \mu^v \} = 2g^{\mu\nu}$$

Sviki preostalih 23 članov (sem prvog)
se mogu svesti na $\mu^0 \mu^1 \mu^2 \mu^3$ i

tako se stize do jednadžbi u postavci
Zadlathy.

7. Pokazati da je

$$[\gamma_5, \sigma^{\mu\nu}] = 0$$

$$\sigma_{\mu\nu} = \frac{i}{2} [\gamma_\mu, \gamma_\nu]$$

$$\begin{aligned} [\gamma_5, \frac{i}{2} [\gamma^\mu, \gamma^\nu]] &= \frac{i}{2} [\gamma_5, \gamma^\mu \gamma^\nu - \gamma^\nu \gamma^\mu] \\ &= \frac{i}{2} (\gamma_5 (\gamma^\mu \gamma^\nu - \gamma^\nu \gamma^\mu) - (\gamma^\mu \gamma^\nu - \gamma^\nu \gamma^\mu) \gamma_5) \\ &= \frac{i}{2} (\gamma_5 \gamma^\mu \gamma^\nu - \gamma_5 \gamma^\nu \gamma^\mu - \gamma^\mu \gamma^\nu \gamma_5 + \gamma^\nu \gamma^\mu \gamma_5) = \end{aligned}$$

$$\boxed{\{\gamma_5, \gamma^\mu\} = 0}$$

$$= \frac{i}{2} (-\gamma^\mu \gamma_5 \gamma^\nu - \gamma_5 \gamma^\nu \gamma^\mu - \gamma^\mu \gamma^\nu \gamma_5 - \gamma^\nu \gamma_5 \gamma^\mu)$$

$$= \frac{i}{2} (\cancel{\gamma^\mu \gamma^\nu \gamma_5} - \cancel{\gamma_5 \gamma^\nu \gamma^\mu} - \cancel{\gamma^\mu \gamma^\nu \gamma_5} + \cancel{\gamma_5 \gamma^\nu \gamma^\mu})$$

$$= 0$$

8. Pokazati da vazi

a) $g_{\mu} g^{\mu} = 4$

b) $g_{\mu} g^{\nu} g^{\mu} = -2 g^{\nu}$

c) $g_{\mu} g^{\alpha} g^{\beta} g^{\mu} = 4 g^{\alpha\beta}$

d) $\epsilon_{\mu\nu} \epsilon^{\mu\nu} = 12$

a) $\{g_{\mu}, g^{\mu}\} = 2 g_{\mu} g^{\mu} = 2 \delta_{\mu}^{\mu} = 8 \Rightarrow g_{\mu} g^{\mu} = 4$
 \hookrightarrow v. pozadi

b) $g_{\mu} g^{\nu} g^{\mu} = -2 g^{\nu}$

$$g_{\mu} g^{\nu} + g^{\nu} g_{\mu} = 2 \delta_{\nu}^{\mu}$$

$$g_{\mu} g^{\nu} = -g^{\nu} g_{\mu} + 2 \delta_{\nu}^{\mu}$$

$$\underline{g_{\mu} g^{\nu} g^{\mu}} = (-g^{\nu} g_{\mu} + 2 \delta_{\nu}^{\mu}) g^{\mu}$$

$$= -g^{\nu} g_{\mu} g^{\mu} + 2 \delta_{\nu}^{\mu} g^{\mu}$$

$$= -4 g^{\nu} + 2 g^{\nu} = \underline{-2 g^{\nu}}$$

c) $g_{\mu} g^{\alpha} g^{\beta} g^{\mu} = 4 g^{\alpha\beta}$

$$\underline{g_{\mu} g^{\alpha} g^{\beta} g^{\mu}} = (-g^{\alpha} g_{\mu} g^{\beta} g^{\mu} + 2 \delta_{\mu}^{\alpha} g^{\beta} g^{\mu}) =$$

$$x^\mu x_\mu = x^0 x^0 + x^1 x^1 + x^2 x^2 + x^3 x^3$$

$$= x^0 x^0 + g_{11} x^1 x^1 + g_{22} x^2 x^2 + g_{33} x^3 x^3$$

$$= (x^0)^2 - (x^1)^2 - (x^2)^2 - (x^3)^2$$

$$= 1 - (-1) - (-1) - (-1)$$

$$= 4$$

$$\begin{aligned}
&= (-\eta^\alpha (-\eta^\beta \eta_\mu + 2\delta_\mu^\beta)) \eta^\mu + 2\eta^\beta \eta^\alpha \\
&= (\eta^\alpha \eta^\beta \eta_\mu \eta^\mu - 2\delta_\mu^\beta \eta^\alpha \eta^\mu + 2\eta^\beta \eta^\alpha) \\
&= (4\eta^\alpha \eta^\beta - 2\eta^\alpha \eta^\beta + 2\eta^\beta \eta^\alpha) \\
&= 2(\eta^\alpha \eta^\beta + \eta^\beta \eta^\alpha) = 2\{\eta^\alpha, \eta^\beta\} \\
&= 4g^{\alpha\beta}
\end{aligned}$$

$$\begin{aligned}
d) \quad \epsilon_{\mu\nu} \epsilon^{\mu\nu} &= \frac{i}{2} [\eta_\mu, \eta_\nu] \frac{i}{2} [\eta^\mu, \eta^\nu] \\
&= -\frac{1}{4} ((\eta_\mu \eta_\nu - \eta_\nu \eta_\mu) (\eta^\mu \eta^\nu - \eta^\nu \eta^\mu)) \\
&= -\frac{1}{4} (\eta_\mu \eta_\nu \eta^\mu \eta^\nu - \eta_\mu \eta_\nu \eta^\nu \eta^\mu - \eta_\nu \eta_\mu \eta^\mu \eta^\nu \\
&\quad + \eta_\nu \eta_\mu \eta^\nu \eta^\mu) \\
&= -\frac{1}{4} (\eta_\mu (-\eta^\mu \eta_\nu + 2\delta_\nu^\mu) \eta^\nu - 16 - 16 + \\
&\quad + \eta_\nu (-\eta^\nu \eta_\mu + 2\delta_\mu^\nu) \eta^\mu) \\
&= -\frac{1}{4} (-\eta_\mu \eta^\mu \eta_\nu \eta^\nu + 2\delta_\nu^\mu \eta_\mu \eta^\nu - 32 + \\
&\quad - \eta_\nu \eta^\nu \eta_\mu \eta^\mu + 2\delta_\mu^\nu \eta_\nu \eta^\mu) \\
&= -\frac{1}{4} (-16 + 8 - 32 - 16 + 8) = 12
\end{aligned}$$

Pokazati da vazi

a) $\text{tr } \mathcal{H}_\mu = 0$

b) $\text{tr}(\mathcal{H}_\mu \mathcal{H}_\nu) = 4 g_{\mu\nu}$

c) $\text{tr}(\mathcal{H}_\mu \mathcal{H}_\nu \mathcal{H}_\rho \mathcal{H}_\sigma) = 4(g_{\mu\nu} g_{\rho\sigma} - g_{\mu\rho} g_{\nu\sigma} + g_{\mu\sigma} g_{\nu\rho})$

d) $\text{tr } \mathcal{H}_5 = 0$

~~e) $\text{tr}(\mathcal{H}_5 \mathcal{H}_\mu \mathcal{H}_\nu) = 0$~~

f) $\text{tr}(\mathcal{H}_5 \mathcal{H}_\mu) = 0$

a) $\{ \mathcal{H}_\mu, \mathcal{H}_5 \} = 0$, $(\mathcal{H}_5)^2 = 1$; $(\mathcal{H}_5)^2 = 1$

$\text{tr } \mathcal{H}_\mu = \text{tr}(\mathcal{H}_\mu \mathcal{H}_5 \mathcal{H}_5) = -\text{tr}(\mathcal{H}_5 \mathcal{H}_\mu \mathcal{H}_5)$
 $= -\text{tr}(\mathcal{H}_5^2 \mathcal{H}_\mu) = -\text{tr}(\mathcal{H}_\mu) \Rightarrow \text{tr } \mathcal{H}_\mu = 0$

b) $\text{tr}(\mathcal{H}_\mu \mathcal{H}_\nu) = ?$

$\{ \mathcal{H}_\mu, \mathcal{H}_\nu \} = 2 g_{\mu\nu}$

matricni Element

$\text{tr}(\mathcal{H}_\mu \mathcal{H}_\nu) = \text{tr}(-\mathcal{H}_\nu \mathcal{H}_\mu + 2 g_{\mu\nu} \mathbf{I})$
 $= -\text{tr}(\mathcal{H}_\mu \mathcal{H}_\nu) + 2 g_{\mu\nu} \text{tr } \mathbf{I} \Rightarrow$

$2 \text{tr}(\mathcal{H}_\mu \mathcal{H}_\nu) = 2 g_{\mu\nu} \cdot 4$
 $= 8 g_{\mu\nu} \Rightarrow \text{tr}(\mathcal{H}_\mu \mathcal{H}_\nu) = 4 g_{\mu\nu}$

$$2) \text{tr} \left[\underbrace{\eta_\mu \eta_\nu \eta_\rho \eta_\sigma}_! \right] =$$

$$= \text{tr} \left[(2g_{\mu\nu} - \eta_\nu \eta_\mu) \eta_\rho \eta_\sigma \right] = 2g_{\mu\nu} \text{tr} \eta_\rho \eta_\sigma -$$

$$- \text{tr} \eta_\nu \eta_\mu \eta_\rho \eta_\sigma = 2g_{\mu\nu} \text{tr} \eta_\rho \eta_\sigma -$$

$$- \text{tr} \eta_\nu (2g_{\mu\rho} - \eta_\rho \eta_\mu) \eta_\sigma = 2g_{\mu\nu} \text{tr} \eta_\rho \eta_\sigma -$$

$$- 2g_{\mu\rho} \text{tr} \eta_\nu \eta_\sigma + \text{tr} \eta_\nu \eta_\rho \eta_\mu \eta_\sigma =$$

$$2g_{\mu\nu} \text{tr} \eta_\rho \eta_\sigma - 2g_{\mu\rho} \text{tr} \eta_\nu \eta_\sigma + \text{tr} \eta_\nu \eta_\rho (2g_{\mu\sigma} - \eta_\sigma \eta_\mu)$$

$$2g_{\mu\nu} \cdot 4g_{\rho\sigma} - 2g_{\mu\rho} \cdot 4g_{\nu\sigma} + 2g_{\mu\sigma} \cdot 4g_{\nu\rho} -$$

$$- \text{tr} \eta_\nu \eta_\rho \eta_\sigma \eta_\mu$$

Dabei

$$\text{tr} \eta_\mu \eta_\nu \eta_\rho \eta_\sigma = 8 (g_{\mu\nu} g_{\rho\sigma} - g_{\mu\rho} g_{\nu\sigma} + g_{\mu\sigma} g_{\nu\rho})$$

$$- \text{tr} \eta_\mu \eta_\nu \eta_\rho \eta_\sigma \Rightarrow$$

$$\text{tr} \eta_\mu \eta_\nu \eta_\rho \eta_\sigma = 4 (g_{\mu\nu} g_{\rho\sigma} - g_{\mu\rho} g_{\nu\sigma} + g_{\mu\sigma} g_{\nu\rho})$$

d) $\text{tr} \eta_5 = 0$

$$\begin{aligned} \text{tr } \beta_5 &= \text{tr} (\beta_5 \beta_0 \beta_0) = -\text{tr} (\beta_0 \beta_5 \beta_0) \\ &= -\text{tr} (\beta_0^2 \beta_5) = -\text{tr} (\beta_5) \Rightarrow \text{tr } \beta_5 = 0 \end{aligned}$$

$$e) \text{tr} (\beta_5 \beta_\mu \beta_\nu) = 0$$

$\beta^\alpha \beta_\alpha = 4$ ← Iz drugog zadatka

$$\begin{aligned} \text{tr} (\beta_5 \beta_\mu \beta_\nu) &= \frac{1}{4} \text{tr} (\beta_5 \beta^\alpha \beta_\alpha \beta_\mu \beta_\nu) \\ &= \frac{1}{4} \text{tr} (\beta_\alpha \beta_\mu \beta_\nu \beta_5 \beta^\alpha) = -\frac{1}{4} \text{tr} (\beta_\alpha \beta_\mu \beta_\nu \beta^\alpha \beta_5) \\ &= -\frac{1}{4} \text{tr} (\beta_5 \beta_\alpha \beta_\mu \beta_\nu \beta^\alpha) = -g_{\mu\nu} \text{tr} (\beta_5) = 0 \end{aligned}$$

Zadatok 8: $4g_{\mu\nu}$

Iskorišćena je i relacija $\{\beta_\mu, \beta^\mu\} = 0$ iz Zad. 7.

$$f) \text{tr} (\beta_5 \beta_\mu) = 0$$

$$\{\beta_5, \beta_\mu\} = 0$$

$$\text{tr} \beta_5 \beta_\mu = -\text{tr} (\beta_\mu \beta_5) = -\text{tr} (\beta_5 \beta_\mu)$$

⇓

$$\text{tr} (\beta_5 \beta_\mu) = 0$$

10. Izračunati

$$\frac{\mu_\mu (1 - \mu_5) (\not{X} - m) \not{e}^\mu}{}$$

$$\left\{ \begin{aligned} \not{k}_\mu \not{k}_5 &= 0 \\ \not{k}_\mu \not{k}_\nu &= 2g_{\mu\nu} \end{aligned} \right.$$

$$\mu_\mu (1 - \mu_5) (\not{X} - m) \not{e}^\mu =$$

$$\mu_\mu (\not{X} - m - \mu_5 \not{X} + m \mu_5) \not{e}^\mu =$$

↑ Prathodni
 ▼ Zadaci

$$\mu_\mu \not{X} \not{e}^\mu - m \underbrace{\mu_\mu \not{e}^\mu}_{\text{circled}} - \mu_\mu \mu_5 \not{X} \not{e}^\mu + m \mu_\mu \mu_5 \not{e}^\mu =$$

$$P_r \mu_\mu \not{e}^\nu \not{e}^\mu - 4m - P_r \mu_\mu \mu_5 \not{e}^\nu \not{e}^\mu - m \mu_5 \underbrace{\mu_\mu \not{e}^\mu}_4 =$$

$$P_r \mu_\mu (-\not{e}^\mu \not{k}^\nu + 2g^{\mu\nu}) - 4m(1 + \mu_5) +$$

$$+ P_r \mu_5 \mu_\mu \not{e}^\nu \not{e}^\mu =$$

$$- 4P_r \not{e}^\nu + 2g^{\mu\nu} P_r \mu_\mu - 4m(1 + \mu_5) + P_r \mu_5 \mu_\mu (-\not{e}^\mu \not{e}^\nu + 2g^{\mu\nu})$$

$$- 4\not{X} + 2\not{X} - 4m(1 + \mu_5) + \underbrace{P_r \mu_5 \mu_\mu \not{e}^\mu \not{e}^\nu}_{\text{circled}} + 2g^{\mu\nu} P_r \mu_5 \mu_\mu =$$

$$- 2\not{X} - 4m(1 + \mu_5) - 4\mu_5 \not{X} + 2\mu_5 \not{X} =$$

$$= -2\not{X} - 4m(1 + \mu_5) - 2\mu_5 \not{X}$$

$$= -2(1 + \mu_5) \not{X} - 4m(1 + \mu_5)$$

$$= -2(1 + \mu_5) (\not{X} + 2m)$$



11. Pokazati da je

$$e^{\eta_5 d} = \cos \sqrt{\eta_5} a^{\mu} + \frac{1}{\sqrt{\eta_5} a^{\mu}} \eta_5 d \sin \sqrt{\eta_5} a^{\mu}$$

$$e^{\eta_5 d} = 1 + \eta_5 d + \frac{1}{2} (\eta_5 d)^2 + \frac{1}{3!} (\eta_5 d)^3 + \dots$$

$$(\eta_5 d)^2 = ?$$

$$(\eta_5 d)^2 = (\eta_5 \eta_{\mu} a^{\mu})^2 = \eta_5 \eta_{\mu} a^{\mu} \eta_5 \eta_{\nu} a^{\nu}$$

$$= \eta_5 \eta_{\mu} \eta_5 \eta_{\nu} a^{\mu} a^{\nu} = -(\eta_5)^2 \eta_{\mu} \eta_{\nu} a^{\mu} a^{\nu}$$

$$= -(\eta_{\mu} \eta_{\nu}) a^{\mu} a^{\nu} = -d^2 \quad \left\{ \eta_5, \eta_{\mu} \right\} = 0$$

$$d^2 = \eta_{\mu} \eta_{\nu} a^{\mu} a^{\nu} = (-\eta_{\nu} \eta_{\mu} + 2g_{\mu\nu}) a^{\mu} a^{\nu}$$

$$= -\eta_{\nu} \eta_{\mu} a^{\mu} a^{\nu} + 2g_{\mu\nu} a^{\mu} a^{\nu}$$

$$= -d^2 + 2g_{\mu\nu} a^{\mu} a^{\nu}$$

$$d^2 = -d^2 + 2a^2$$

$$\Downarrow$$
$$d^2 = a^2$$

$$(\eta_5 d)^2 = -a^2$$

$$(k_5 d)^3 = (k_5 d)^2 (k_5 d) = -a^2 (k_5 d)$$

$$(k_5 d)^4 = (k_5 d)^2 (k_5 d)^2 = a^4$$

$$(k_5 d)^5 = (k_5 d)^4 (k_5 d) = a^4 (k_5 d)$$

⋮

$$e^{k_5 d} = \left(1 + k_5 d - \frac{1}{2} a^2 + \frac{a^2}{3!} (k_5 d) + \frac{1}{4!} a^4 + \frac{a^4}{5!} (k_5 d) + \dots \right)$$

$$= \left(1 - \frac{1}{2!} a^2 + \frac{a^4}{4!} + \dots \right) + (k_5 d) \left(1 - \frac{a^2}{3!} + \frac{a^4}{5!} + \dots \right)$$

$$= \cos \sqrt{a^2} + \frac{1}{\sqrt{a^2}} (k_5 d) \sin \sqrt{a^2}$$

$$a^2 = a_n a^m$$

$$e^{ix} = \sum_{n=0}^{\infty} \frac{(ix)^n}{n!} = 1 + ix - \frac{x^2}{2!} - \frac{ix^3}{3!} + \frac{x^4}{4!} + \frac{ix^5}{5!} - \frac{x^6}{6!} + \dots$$

$$\left(1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \frac{x^6}{6!} + \dots \right) + i \left(x - \frac{x^3}{3!} + \frac{x^5}{5!} + \dots \right)$$

$$e^{ix} = \cos x + i \sin x \quad (\text{Euler})$$

$$\cos x = 1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \frac{x^6}{6!} + \dots$$

$$\sin x = x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots$$

12. Neka je data skup matrica

$$\Gamma^a = \{ I, \gamma^m, \gamma^5, \gamma^m \gamma^5, \gamma^{\mu\nu} \}$$

sa osobinama

1. $(\Gamma^a)^2 = I$

2. $\Gamma^a \Gamma^b = \pm \Gamma^c \quad \begin{cases} a=b, & \pm=1 \\ a \neq b, & \pm i, \mp i \end{cases}$

3. Ako je $\Gamma^a \neq I$, onda uvek postoji neka Γ^b takvo da vazi

$$\Gamma^b \Gamma^a \Gamma^b = -\Gamma^a \Leftrightarrow \{ \Gamma^b, \Gamma^a \} = 0$$

4. Ako je $\Gamma^a \neq I$

$$\text{tr } \Gamma^a = 0, \quad a = 2, \dots, 16$$

Pokazati da Γ^a -ovi cine LNZ skup i da se proizvoljna matrica $A \in C^{4 \times 4}$ moze zapisati

kao

$$A = \sum_a \Theta_a \Gamma^a$$

gde je $\Theta_a = \frac{1}{4} \text{tr}(A \Gamma^a)$

LNZ $\sum_{a=1}^{16} \beta_a \Gamma^a$

$$\sum_{a=1}^{16} \beta_a \Gamma^a = 0$$

$\beta_a = 0, \forall a \Leftrightarrow \Gamma^a \text{ LNZ}$

$$\text{tr} \sum_{a=1}^{16} \beta_a \Gamma^a = 0$$

$$\text{tr} \beta_1 \Gamma^1 + \text{tr} \sum_{a=2}^{16} \beta_a \Gamma^a = 0$$

$$4\beta_1 + \sum_{a=2}^{16} \beta_a \cdot 0 = 0$$

$$\boxed{\beta_1 = 0}$$

$$\sum_{a=2}^{16} \beta_a \Gamma^a = 0 \quad / \Gamma^b$$

$$\sum_{a=2}^{16} \beta_a \Gamma^a \Gamma^b = 0$$

$$\beta_a (\Gamma^a)^2 + \sum_{a \neq b} \beta_a \Gamma^a \Gamma^b = 0$$

1

$$\beta_a + \sum_{a \neq b} \beta_a \Gamma^a \Gamma^b = 0 \quad / \text{tr}$$

Вспомогательная (вспомогательная) 4×4

$$4\beta_a + \sum_{a \neq b} \beta_a \operatorname{tr}(\Gamma^a \Gamma^b) = 0$$

Γ_3 .

$$\{\Gamma^b, \Gamma^a\} = 0 \Rightarrow \operatorname{tr} \Gamma^a \Gamma^b = 0$$

$$\beta_a = 0, \quad a = 2, \dots, 16$$

$$A = \sum_a \theta_a \Gamma^a$$

$$A \Gamma^b = \sum_a \theta_a \Gamma^a \Gamma^b$$

$$= \theta_b I + \sum_{a \neq b} \theta_a \Gamma^a \Gamma^b$$

$$\operatorname{tr} A \Gamma^b = 4\theta_b + \sum_{a \neq b} \theta_a \operatorname{tr}(\Gamma^a \Gamma^b)$$

$$\boxed{\theta_b = \frac{1}{4} \operatorname{tr} A \Gamma^b}$$

13. Pokazati da je za Dirakovu $\bar{\psi}$ -nu četvorovektor struje dat izrazom

$$j^\mu = \bar{\psi} \gamma^\mu \psi$$

gde je $\bar{\psi} = \psi^\dagger \gamma^0$

$\bar{\psi}$ - adjoint

(ψ надыучето, српски) (1)

$$(i \gamma^\mu \partial_\mu - m) \psi = 0$$

$$(\gamma^\mu)^\dagger = \gamma^0 \gamma^\mu \gamma^0 \Rightarrow (\gamma^\mu)^\dagger \gamma^0 \stackrel{(\gamma^0)^2 = 1}{=} \gamma^0 \gamma^\mu$$

13. (1), ~~kompleksna~~ ~~konjugovan~~ (adjungovan \Rightarrow)

$$\psi^\dagger (-i (\gamma^\mu)^\dagger \partial_\mu - m) = 0 \quad | \gamma^0$$

$$\psi^\dagger (-i (\gamma^\mu)^\dagger \gamma^0 \partial_\mu - m \gamma^0) = 0$$

$$\psi^\dagger \gamma^0 (-i \gamma^\mu \partial_\mu - m) = 0$$

$$\bar{\psi} (-i \gamma^\mu \partial_\mu - m) = 0$$

$$\bar{\psi} (i \gamma^\mu \overleftarrow{\partial}_\mu + m) = 0 \quad (2)$$

Pomnoziti poslednju $\bar{\psi}$ -m sa ψ s desna
a (1) sa $\bar{\psi}$ s leva:

$$\bar{\psi} (-i \gamma^\mu \overleftarrow{\partial}_\mu - m) \psi = 0$$

$$\bar{\psi} (i \gamma^\mu \overrightarrow{\partial}_\mu - m) \psi = 0$$

} \Rightarrow

$$i \bar{\psi} \gamma^\mu \overleftarrow{\partial}_\mu \psi = -m \bar{\psi} \psi$$

$$i \bar{\psi} \gamma^\mu \overrightarrow{\partial}_\mu \psi = m \bar{\psi} \psi$$

\oplus

$$i \partial_\mu (\bar{\psi} \gamma^\mu \psi) = 0$$

$\Downarrow \partial_\mu j^\mu = 0$ - jednačina kontinuiteta

$$j^\mu = \bar{\psi} \gamma^\mu \psi$$

14. Rešiti Dirakovu \hat{H} -nu za slobodnu česticu, pretpostavljajući da Hamiltonijan ne zavisi eksplicitno od vremena.

Dirakova \hat{H} -na \hat{H}

$$(i \gamma^\mu \partial_\mu - m) \Psi(x) = 0 \quad ; \quad x = x^\mu$$

$\Psi(x)$ - četvorkomponentni spinor

$$\Psi(x) = \begin{pmatrix} \psi_1(x) \\ \psi_2(x) \\ \psi_3(x) \\ \psi_4(x) \end{pmatrix}$$

Hamiltonijan ne zavisi eksplicitno od vremena pa se rešenje može tražiti u obliku

$$\Psi(x) = \Psi(\vec{r}, t) = \phi(\vec{r}) e^{-\frac{i}{\hbar} E t} \quad (1)$$

sto znači da se Dirakova \hat{H} -na svodi na

$$\hat{H}_0 \phi(\vec{r}) = E \phi(\vec{r}) \quad \left(\begin{array}{l} \text{stacionarna} \\ \hat{H}\text{-na} \end{array} \right) \text{ uveriti se}$$

Budući da je zadovoljeno $[\hat{H}_0, \vec{P}] = 0$, po analogiji sa kvantnom mehanikom sledi

$$\phi_p(\vec{r}) = U_p e^{\frac{i}{\hbar} \vec{P} \cdot \vec{r}} \quad (2)$$

odnosno

$$\Psi(x) = U_p e^{\frac{i}{\hbar} \vec{P} \cdot \vec{r}} e^{-\frac{i}{\hbar} E t} = U_p e^{\frac{i}{\hbar} (\vec{r} \cdot \vec{P} - E t)}$$

U_p - konstantan spinor

Ovo partikularno rešenje mora da zadovolji Dirakovu granu

$$(\not{p} - m) \psi(x) = 0$$

odnosno

$$(\not{p} - m) U_p = 0 \quad (3)$$

U_p zapisujemo u obliku

$$U_p = \begin{pmatrix} \psi \\ \chi \end{pmatrix}_{4 \times 1} \quad (4)$$

$$\not{p} = p_\mu \gamma^\mu = p_0 \gamma^0 - \vec{p} \cdot \vec{\gamma}$$

$$= p_0 \begin{pmatrix} \mathbb{I} & 0 \\ 0 & -\mathbb{I} \end{pmatrix} - \vec{p} \begin{pmatrix} 0 & \vec{\sigma} \\ -\vec{\sigma} & 0 \end{pmatrix}$$

$$= \begin{pmatrix} E & 0 \\ 0 & -E \end{pmatrix} - \begin{pmatrix} 0 & \vec{p} \cdot \vec{\sigma} \\ -\vec{p} \cdot \vec{\sigma} & 0 \end{pmatrix}$$

$$= \begin{pmatrix} E & -\vec{\sigma} \cdot \vec{p} \\ \vec{\sigma} \cdot \vec{p} & -E \end{pmatrix}_{4 \times 4}$$

↓

Блок-матрица

Koristeći (3) i (4), sledi

$$\left[\begin{pmatrix} E & -\vec{b} \cdot \vec{p} \\ \vec{b} \cdot \vec{p} & -E \end{pmatrix} - \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix}_{4 \times 4} \right] \begin{pmatrix} \varphi \\ \chi \end{pmatrix}_{4 \times 1} = 0$$

$$\begin{pmatrix} E-m & -\vec{b} \cdot \vec{p} \\ \vec{b} \cdot \vec{p} & -E-m \end{pmatrix} \begin{pmatrix} \varphi \\ \chi \end{pmatrix} = 0 \quad (5)$$

Homogeni sistem (5) ima rešenje ako je determinanta odgovarajuće matrice nula

$$-(E-m)^2 + (\vec{b} \cdot \vec{p})^2 = 0$$

$$-E^2 + m^2 + \vec{p} \cdot \vec{p} + i \vec{b} \cdot (\vec{p} \times \vec{p}) = 0$$

$$-E^2 = -m^2 - p^2$$

$$E^2 = m^2 + p^2$$

⌈ Krantna mehanika

$$(\vec{b} \cdot \vec{a})(\vec{b} \cdot \vec{b}) = \vec{a} \cdot \vec{b} + i \vec{b} \cdot (\vec{a} \times \vec{b})$$

$$E = \pm \sqrt{p^2 + m^2} = \pm E_p$$

⌈ (5) ⇒

$$(E-m)\varphi - \vec{b} \cdot \vec{p} \chi = 0$$

$$(6) \text{ za } E = -E_p$$

$$\vec{b} \cdot \vec{p} \varphi - (E+m)\chi = 0$$

$$(7) \text{ za } E = E_p$$

$$I_z (7), E = E_p$$

$$\chi = \frac{\vec{\sigma} \cdot \vec{p}}{E_{ptm}} \psi$$

Spinor onda glasi

$$U_p = \begin{pmatrix} \psi \\ \frac{\vec{\sigma} \cdot \vec{p}}{E_{ptm}} \psi \end{pmatrix} \equiv U(E_p, \vec{p})$$

$$I_z (6), E = -E_p$$

$$\psi = - \frac{\vec{\sigma} \cdot \vec{p}}{E_{ptm}} \chi$$

Spinor onda glasi

$$U_p = \begin{pmatrix} \psi \\ \chi \end{pmatrix} = \begin{pmatrix} - \frac{\vec{\sigma} \cdot \vec{p}}{E_{ptm}} \chi \\ \chi \end{pmatrix} \equiv U(-E_p, \vec{p})$$

Ako se uvedu nove oznake

$$v(\vec{p}) = U(-E_p, -\vec{p}) = \begin{pmatrix} \frac{\vec{\sigma} \cdot \vec{p}}{E_{ptm}} \chi \\ \chi \end{pmatrix}$$

$$i u(\vec{p}) = U(E_p, \vec{p}) = \begin{pmatrix} \psi \\ \frac{\vec{\sigma} \cdot \vec{p}}{E_{ptm}} \psi \end{pmatrix}$$

Kiž rešenja Diranove jne glase

$$u(\vec{p}) e^{-ip \cdot x} \quad i \quad v(\vec{p}) e^{ip \cdot x}$$

Iz prethodnog zadatka je poznato

$$j^\mu = \bar{\psi} \gamma^\mu \psi$$

odnosno

$$j^0 = \rho = \psi^\dagger \underbrace{\gamma^0 \gamma^0}_{1} \psi = \psi^\dagger \psi \quad \left(\rho = |\psi|^2 \right)$$

Posmatrajmo, formalno, ρ kao statistički operator u prostoru Hilbertovom, pri čemu važi

$$\text{tr} \rho = 1$$

Na osnovi ovog se mogu normalizovati dobijena LNE rešenja.

$$\psi = N U(\vec{p}) e^{-i\vec{p} \cdot \vec{x}}$$

$$\rho = |N|^2 U^\dagger(\vec{p}) U(\vec{p})$$

$$= |N|^2 \left(\psi^\dagger \quad \psi^\dagger \frac{(\vec{\sigma} \cdot \vec{p})^\dagger}{E_{ptm}} \right) \begin{pmatrix} \psi \\ \frac{\vec{\sigma} \cdot \vec{p}}{E_{ptm}} \psi \end{pmatrix}$$

$$= |N|^2 \psi^\dagger \left(1 \quad \frac{\vec{\sigma} \cdot \vec{p}}{E_{ptm}} \right) \begin{pmatrix} 1 \\ \frac{\vec{\sigma} \cdot \vec{p}}{E_{ptm}} \end{pmatrix} \psi ; \quad \psi = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

$$= |N|^2 \left(1 + \frac{(\vec{\sigma} \cdot \vec{p})^2}{(E_{ptm})^2} \right) \underbrace{\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}}_1 \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

✓
↓
uslov

Vec je gradeno $(\vec{0} \cdot \vec{p})^2 = E_p^2 - m^2$

$$S = |N|^2 \left(1 + \frac{E_p^2 - m^2}{(E_{p+m})^2} \right) =$$

$$= |N|^2 \left(1 + \frac{E_p - m}{E_{p+m}} \right) = |N|^2 \frac{E_{p+m} + E_p - m}{E_{p+m}}$$

$$= |N|^2 \frac{2E_p}{E_{p+m}}$$

ta $S = 1 \Rightarrow N = \sqrt{\frac{E_{p+m}}{2E_p}} e^{i\phi}$

Dakle, jedno rešenje je

$$\psi(x) = \sqrt{\frac{E_{p+m}}{2E_p}} u(\vec{p}) e^{-i\vec{p} \cdot x}$$

$$= \frac{1}{\sqrt{2E_p}} u(\vec{p}) e^{-i\vec{p} \cdot x}$$

Opšte rešenje

$$\psi(x) = \frac{1}{(2\pi)^{3/2}} \sum_{r=1}^2 \int d^3\vec{p} (c_r(\vec{p}) \psi_r^+(x) + d_r^\dagger(\vec{p}) \psi_r^-(x))$$

$$= \frac{1}{(2\pi)^{3/2}} \sum_{r=1}^2 \int \frac{d^3\vec{p}}{\sqrt{2E_p}} (c_r(\vec{p}) u_r(\vec{p}) e^{-i\vec{p} \cdot x} + d_r^\dagger(\vec{p}) v_r(\vec{p}) e^{i\vec{p} \cdot x})$$

15. Pokazati nam izgleda rešenje Dirakove \hat{H} -n ψ za $E > 0$ u nerelativističkom (niznom).

$$\psi = N \begin{pmatrix} \varphi \\ \frac{\vec{\sigma} \cdot \vec{p}}{E + m_0 c^2} \varphi \end{pmatrix} e^{-i\vec{p} \cdot \vec{x}}$$

$$\chi = \frac{\vec{\sigma} \cdot \vec{p}}{E + m_0 c^2} \varphi$$

$$\begin{aligned} E_p^2 &= \vec{p}^2 c^2 + m_0^2 c^4 \\ &= m_0^2 c^4 \left(1 + \frac{\vec{p}^2}{m_0^2 c^2} \right) \end{aligned}$$

$$E_p = m_0 c^2 \sqrt{1 + \frac{\vec{p}^2}{m_0^2 c^2}}$$

$$\approx m_0 c^2 \left(1 + \frac{1}{2} \frac{\vec{p}^2}{m_0^2 c^2} \right)$$

$$\left[(1+x)^{\frac{1}{2}} \approx 1 + \frac{1}{2} x \right]$$

$$x \ll 1$$

$$E = m_0 c^2 + \frac{\vec{p}^2}{2m_0}$$

$$\Rightarrow \boxed{E_p - m_0 c^2 = \frac{\vec{p}^2}{2m_0}}$$

$$E_p + m_0 c^2 - m_0 c^2 - m_0 c^2 = \frac{\vec{p}^2}{2m_0}$$

$$\boxed{E_p + m_0 c^2 = 2m_0 c^2 + \frac{\vec{p}^2}{2m_0}}$$

$$m_0 \gg \frac{\vec{p}^2}{2m_0}$$

$$\chi = \frac{\vec{\sigma} \cdot \vec{p}}{E_p + m_0} \psi \Rightarrow \left(\vec{\sigma} \cdot \vec{p} = \sqrt{E_p^2 - m_0^2} \right)$$

$$\frac{\chi}{\psi} = \frac{\sqrt{E_p^2 - m_0^2}}{E_p + m_0} = \sqrt{\frac{E_p^2 - m_0^2}{(E_p + m_0)^2}} = \sqrt{\frac{E_p - m_0}{E_p + m_0}} \approx$$

$$= \sqrt{\frac{\frac{p^2}{2m_0}}{2m_0 + \frac{p^2}{2m_0}}} = \sqrt{\frac{1}{\frac{2m_0}{p^2} + 1}} = \sqrt{\frac{1}{\frac{4m_0^2}{p^2} + 1}} \ll 1$$

Onda je

$$\Psi = N \begin{pmatrix} \varphi \\ 0 \end{pmatrix} e^{-ipx}$$

$$\varphi = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \text{ ili } \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

pa će biti

$$\Psi = N \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} e^{-ipx}$$

ili

$$\Psi = N \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} e^{-ipx}$$

rešenje u
nerelativističkom
limesu

110. Pona zati vazeenje pelacovca

$$\sum_{r=1}^2 U_r(\vec{p}) \bar{U}_r(\vec{p}) = \not{p} + m$$

$$\sum_{r=1}^2 U_r(\vec{p}) \bar{U}_r(\vec{p}) = \not{p} - m$$

$$U_r(\vec{p}) = \sqrt{E_{p+m}} \begin{pmatrix} \varphi_r \\ \frac{\vec{\sigma} \cdot \vec{p}}{E_{p+m}} \varphi_r \end{pmatrix}$$

$$\sum_{r=1}^2 U_r(\vec{p}) \bar{U}_r(\vec{p}) = \sum_{r=1}^2 U_r(\vec{p}) U_r^\dagger(\vec{p}) \not{p}^0 =$$

$$= U_1(\vec{p}) U_1^\dagger(\vec{p}) \not{p}^0 + U_2(\vec{p}) U_2^\dagger(\vec{p}) \not{p}^0$$

$$= (E_{p+m}) \left[\begin{pmatrix} \varphi_1 \\ \frac{\vec{\sigma} \cdot \vec{p}}{E_{p+m}} \varphi_1 \end{pmatrix} \begin{pmatrix} \varphi_1^\dagger & \frac{\vec{\sigma} \cdot \vec{p}}{E_{p+m}} \varphi_1^\dagger \end{pmatrix} \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix} + \right.$$

$$\left. \begin{pmatrix} \varphi_2 \\ \frac{\vec{\sigma} \cdot \vec{p}}{E_{p+m}} \varphi_2 \end{pmatrix} \begin{pmatrix} \varphi_2^\dagger & \frac{\vec{\sigma} \cdot \vec{p}}{E_{p+m}} \varphi_2^\dagger \end{pmatrix} \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix} \right] = \begin{cases} \varphi_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ \varphi_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \end{cases}$$

$$\begin{pmatrix} E_{p+m} & -\vec{\sigma} \cdot \vec{p} \\ \vec{\sigma} \cdot \vec{p} & -(E_{p+m}) \end{pmatrix} = \not{p} + m$$

(v. zadatok 14)

Drugo za domaći!